

§29. The Flierl-Petviashvili Equation and the Zonal Flow

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The drift instability in tokamak plasma exhibits distinct types of structures on different spacial scales. A multiple space-time scale analysis of the governing equations (continuity, ion momentum and adiabaticity) has revealed that on scales intermediate between the large linear cells and the sonic Larmor radius (ρ_s) the stationary states are solutions of the Flierl-Petviashvili equation

$$\Delta\phi = \alpha\phi - \beta\phi^2 \quad (1)$$

$\alpha = (1 - v_*/u)/\rho_s^2$, $\beta = (c_s^2/2u^2)d(L_n^{-1})/dx$. Starting from the one-dimensional solution $\phi(y) = 3\alpha/(2\beta)\cosh^{-2}(\sqrt{\alpha}y/2)$ we look for the dependence of the complex- y singularities on the other coordinate, x . We find that the solution is a doubly periodic meromorphic function

$$\phi_s(x, y) = \frac{\alpha}{2\beta} + s\wp(ia y + ibx + \omega|g_2, g_3) \quad (2)$$

$g_2 = 3\alpha^2/(s\beta)^2$ and g_3 are the invariants of the Weierstrass elliptic functions, \wp . Eq.(2) is the exact, periodic, solution of the Flierl-Petviashvili equation. It has a one dimensional geometry consisting of layers of periodic flow (Fig.1) with an orientation in the (y, x) plane given by a/b which verify $a^2 + b^2 = s\beta/6$. g_3 and s are constants of integration fixed by the physical conditions when the system evolves to this solution. When $g_3 = 0$ the width of the layer of poloidal flow is $\delta x = 3.924K(k')(1 - v_*/u)^{-1/2}\rho_s$.

The pattern is the same as the zonal flow in plasma. Experimental studies on Doublet III-D tokamak have obtained $\tilde{\varphi}_{rms} > 10V$, $\lambda_r \in (15...30)\rho_s$, and a flow shear $\omega_{E \times B} \sim 2 \times 10^5 s^{-1}$. From Eq.(2) we obtain: $\tilde{\varphi}_{rms} > 17V$, $\lambda_r \simeq 17.4\rho_s$, $\omega_{E \times B} \sim 2.2 \times 10^5 s^{-1}$. Results of gyrokinetic simulations at Lausanne have obtained $\delta x \sim [8...13]\rho_s$, $E_r \in [-24, +24] \times 10^3 V/m$ and the flow shear rate $\omega_{E \times B} \in [-5, +5] \times 10^{-3} \Omega_i$. We obtain $\delta x \simeq 13\rho_s$, $E_r \in [-23, +23] \times 10^3 V/m$ and $\omega_{E \times B} \in [-10, +10] \times 10^{-3} \Omega_i$.

In the presence of a perturbation the flow evolves in time. The scalar nonlinearity equation can

be written in the form ($4\eta^2$ is a constant)

$$\frac{\partial}{\partial t} (1 - \nabla_\perp^2) \phi = \frac{\partial}{\partial y} [(-\nabla_\perp^2 + 4\eta^2) \phi] - \phi \frac{\partial \phi}{\partial y}$$

We obtain the linear dispersion relation

$$\frac{1}{k_y} = \frac{\omega \pm Q^{1/2}}{P + 1} \left\{ 1 \mp \left[1 + \frac{2(P + 1)}{(-\omega + Q^{1/2})^2} \right]^{1/2} \right\}$$

where $P \equiv -4\eta^2 + (\phi_s|_{\min} + \phi_s|_{\max})/2$ and $Q \equiv (\delta x)^2 (\phi_s|_{\max} - \phi_s|_{\min})^2 \equiv (\delta x)^2 (\delta \phi_s)^2$. The vortical structures have $\lambda_y \sim 5\rho_s$.

Extensive numerical simulations have been done for studying the stability and the structural stability (*i.e.* including both the scalar and the vectorial nonlinearities) of ϕ_s . For small monopolar perturbations the total flow is stable for more than $12 \times 10^3 \Omega_i$. The monopole is reshaped at early stages for equality of the tangential flows and then is stably advected ($v_*/u \lesssim 1$). This result may be relevant for the stability of the Red Spot, the long-lived vortex embedded in the zonal flows of the atmosphere of Jupiter, a geometry very similar to ours. A small amplitude dipolar initial perturbation is also stable for similar durations.

This exact solution has the necessary characteristics to represent the zonal flow as a coherent nonlinear structure and strongly supports the role of the scalar nonlinearity in the phenomenology of the zonal flows [1].

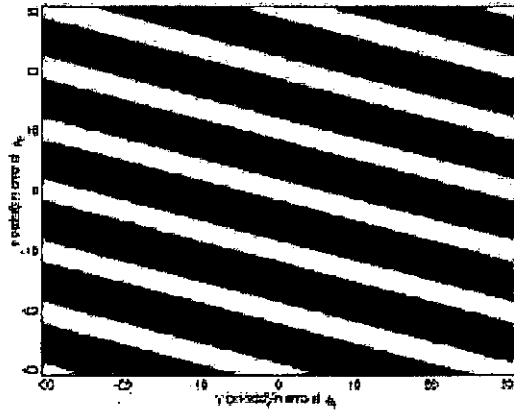


Figure 1: Amplitude of the perturbed potential from the solution ϕ_s for $a/b = 0.9$

References

- [1] F. Spineanu, M. Vlad, K. Itoh, H. Sanuki and S.-I. Itoh, Phys. Rev. Letters, (2004) in print.